

Communication protocols

Quantum communication

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October 2, 2020



Communication protocols

Review of formalism

Superdense coding

Quantum teleportation

Recall: qubit states

Quantum states of a 2-level quantum system (*qubit*):

$$|\phi\rangle = \alpha |0\rangle + \beta |1\rangle \in \mathcal{H}, \quad \alpha, \beta \in \mathbb{C}.$$

Hermitian product: if $|\psi\rangle = \delta |0\rangle + \varepsilon |1\rangle$,

$$\begin{aligned} \langle \phi | \psi \rangle &= \alpha^* \delta + \beta^* \varepsilon \\ &= [\alpha \ \beta]^* \begin{bmatrix} \delta \\ \varepsilon \end{bmatrix} \\ &= |\phi\rangle^\dagger |\psi\rangle \end{aligned}$$

Measurement

When $|\phi\rangle = \alpha|0\rangle + \beta|1\rangle$ is measured:

$|\phi\rangle$ projects to a (random) classical state $\mathcal{M}|\phi\rangle \in \{|0\rangle, |1\rangle\}$ with

$$\mathbb{P}[\mathcal{M}|\phi\rangle = |0\rangle] = \frac{|\alpha|^2}{|\alpha|^2 + |\beta|^2}, \quad \mathbb{P}[\mathcal{M}|\phi\rangle = |1\rangle] = \frac{|\beta|^2}{|\alpha|^2 + |\beta|^2}.$$

In particular: if $|\phi\rangle$ is *normalized*, this reduces to

$$\mathbb{P}[\mathcal{M}|\phi\rangle = |0\rangle] = |\alpha|^2, \quad \mathbb{P}[\mathcal{M}|\phi\rangle = |1\rangle] = |\beta|^2.$$

More generally

$|\phi\rangle$ can be measured with respect to any orthonormal basis $|\varphi_1\rangle, |\varphi_2\rangle$

$$|\phi\rangle = \langle\varphi_1|\phi\rangle|\varphi_1\rangle + \langle\varphi_2|\phi\rangle|\varphi_2\rangle$$

$$\implies \mathbb{P}[\mathcal{M}|\phi\rangle = |\varphi_i\rangle] = |\langle\varphi_i|\phi\rangle|^2$$

Example:

$$|\phi^+\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}} \quad |\phi^-\rangle = \frac{|0\rangle - |1\rangle}{\sqrt{2}} \quad \text{orthonormal basis}$$

2-qubit system

Orthonormal basis of the 4-dimensional space $\mathcal{H} \otimes \mathcal{H}$:

$$|0\rangle \otimes |0\rangle, \quad |0\rangle \otimes |1\rangle, \quad |1\rangle \otimes |0\rangle, \quad |1\rangle \otimes |1\rangle$$

But also the 4 Bell states:

$$\begin{aligned} |\Phi^+\rangle &= \frac{|0\rangle \otimes |0\rangle + |1\rangle \otimes |1\rangle}{\sqrt{2}}, & |\Phi^-\rangle &= \frac{|0\rangle \otimes |0\rangle - |1\rangle \otimes |1\rangle}{\sqrt{2}} \\ |\Psi^+\rangle &= \frac{|0\rangle \otimes |1\rangle + |1\rangle \otimes |0\rangle}{\sqrt{2}}, & |\Psi^-\rangle &= \frac{|0\rangle \otimes |1\rangle - |1\rangle \otimes |0\rangle}{\sqrt{2}} \end{aligned}$$

Entanglement

Recall that a state $|\Psi\rangle \in \mathcal{H} \otimes \mathcal{H}$ is *entangled* if it cannot be written as a simple combination of two single qubit states:

$$\forall_{|\varphi\rangle, |\psi\rangle \in \mathcal{H}} \quad |\Psi\rangle = |\varphi\rangle \otimes |\psi\rangle.$$

Exercise: is the following state entangled?

$$\frac{1}{2}(|0\rangle \otimes |0\rangle - |0\rangle \otimes |1\rangle + |1\rangle \otimes |0\rangle - |1\rangle \otimes |1\rangle)$$

No-cloning theorem

Theorem: *Quantum information cannot be copied.*

Precise statement: there exist no unitary operator $C : \mathcal{H} \rightarrow \mathcal{H} \otimes \mathcal{H}$ such that

$$\forall |\phi\rangle \in \mathcal{H} \quad C|\phi\rangle = |\phi\rangle \otimes |\phi\rangle.$$

...of course! But even if you start with a blank qubit:

there is no unitary operator $C : \mathcal{H} \otimes \mathcal{H} \rightarrow \mathcal{H} \otimes \mathcal{H}$ such that

$$\forall |\phi\rangle \in \mathcal{H} \quad C(|\phi\rangle \otimes |0\rangle) = |\phi\rangle \otimes |\phi\rangle.$$

Reason

Suppose

$$\begin{cases} C(|0\rangle \otimes |0\rangle) = |0\rangle \otimes |0\rangle \\ C(|1\rangle \otimes |0\rangle) = |1\rangle \otimes |1\rangle \end{cases}$$

Then what is

$$C((|0\rangle + |1\rangle) \otimes |0\rangle) ?$$

A cloning operator cannot respect superposition of quantum states.

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Superdense coding

- Bennett & Weisner (1992)
- **Two** classical bits can be transmitted on **one** quantum bit
- An entangled pair is used (and needed)
- A form of secure quantum communication
- Experimentally tested in various settings

Superdense coding: the protocol

Alice has two classical bits b_1 and b_2 that she wants to send to Bob.



$b_1 b_2$



Superdense coding: the protocol

1) Preparation (can be done in advance)



Charlie prepares an entangled pair of qubits in Bell state

$$|\Phi^+\rangle = \frac{|0\rangle_A \otimes |0\rangle_B + |1\rangle_A \otimes |1\rangle_B}{\sqrt{2}}.$$

2) Sharing (can be done in advance)

Charlie sends qubit A to Alice, qubit B to Bob.

Superdense coding: the protocol

3) Encoding



Alice modifies her qubit according to the two classical bits $b_1 b_2$ she wants to encode:

- if $b_1 b_2 = 00$ she does nothing, i.e. applies $I = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$ to A

$$|\Phi^+\rangle = \frac{|0\rangle_A \otimes |0\rangle_B + |1\rangle_A \otimes |1\rangle_B}{\sqrt{2}}$$

Superdense coding: the protocol

- if $b_1 b_2 = 01$ she flips the phase of her qubit, i.e. applies $Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$ to A

$$|\Phi^-\rangle = \frac{|0\rangle_A \otimes |0\rangle_B - |1\rangle_A \otimes |1\rangle_B}{\sqrt{2}}$$

- if $b_1 b_2 = 10$ she negates her qubit, i.e. applies $X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$ to A

$$|\Psi^+\rangle = \frac{|0\rangle_A \otimes |1\rangle_B + |1\rangle_A \otimes |0\rangle_B}{\sqrt{2}}$$

- if $b_1 b_2 = 11$ she applies both X and Z (order doesn't really matter: why?)

$$|\Psi^-\rangle = \frac{|0\rangle_A \otimes |1\rangle_B - |1\rangle_A \otimes |0\rangle_B}{\sqrt{2}}$$

Superdense coding: the protocol

Summary of encoding: $b_1 b_2$ encoded as $Z^{b_2} X^{b_1} |\Phi^+\rangle$



$$00 \mapsto |\Phi^+\rangle$$

$$01 \mapsto |\Phi^-\rangle$$

$$10 \mapsto |\Psi^+\rangle$$

$$11 \mapsto |\Psi^-\rangle$$

Superdense coding: the protocol

4) Sending

Alice sends her qubit to Bob through a quantum channel; now Bob has the full entangled pair.



5) Decoding

Making a measurement in the orthonormal basis $|\Phi^+\rangle, |\Phi^-\rangle, |\Psi^+\rangle, |\Psi^-\rangle$ Bob is able to recover the initial pair of bits.

Superdense coding: the protocol

Reformulation: let T denote the unitary transformation for which

$$|0\rangle \otimes |0\rangle \mapsto |\Phi^+\rangle, \quad |0\rangle \otimes |1\rangle \mapsto |\Phi^-\rangle, \quad |1\rangle \otimes |0\rangle \mapsto |\Psi^+\rangle, \quad |1\rangle \otimes |1\rangle \mapsto |\Psi^-\rangle$$

$$T = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 1 & -1 \\ 1 & -1 & 0 & 0 \end{bmatrix}$$

Then Bob can just perform T^\dagger on the qubit pair to recover $|b_1\rangle \otimes |b_2\rangle$.

Superdense coding: discussion

- A compression factor of 2 can be achieved... (at a cost!)
- EPR pairs can be used to "store bandwidth" for future use
- If Eve measures Alice's qubit in transit: she will randomly get $|0\rangle$ or $|1\rangle$ no matter what the transmitted bits were
- Unfortunately Bob would then receive either a

$$|0\rangle \otimes |0\rangle, \quad |1\rangle \otimes |1\rangle, \quad |0\rangle \otimes |1\rangle \quad \text{or} \quad |1\rangle \otimes |0\rangle :$$

b_1 survives but b_2 is lost (gets random output)

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Quantum teleportation



Quantum teleportation



Quantum teleportation

- Bennett, Brassard, Crépeau, Jozsa, Peres, Wootters (1993)
- In superdense coding: 2 classical bits are sent using 1 quantum bit
- Quantum teleportation can be thought of as the reverse:
1 quantum bit is sent using 2 classical bits of information
- No quantum system is physically transported: only its *state* is
- So: quantum information can be sent through a classical channel...
provided both parties share an entangled pair.

Quantum teleportation: the protocol

Alice has a qubit in state $|\phi\rangle_A = \alpha|0\rangle_A + \beta|1\rangle_A$ that she wants to send to Bob.



1) Preparation (can be done in advance)

Charlie prepares an entangled pair of qubits in Bell state

$$|\Phi^+\rangle = \frac{|0\rangle_C \otimes |0\rangle_B + |1\rangle_C \otimes |1\rangle_B}{\sqrt{2}}.$$

Superdense coding: the protocol

2) **Sharing** (can be done in advance)



Charlie gives qubit C to Alice, qubit B to Bob.

Now Alice has two qubits A and C , Bob has qubit B ; total state of the system is

$$|\phi\rangle_A \otimes |\Phi^+\rangle_{CB} = \frac{\alpha |000\rangle + \alpha |011\rangle + \beta |100\rangle + \beta |111\rangle}{\sqrt{2}}.$$

Superdense coding: the protocol

3) Measurement



Alice measures her two qubits in the Bell basis; the system degenerates equiprobably into one of the 4 states:

$$|\Phi^+\rangle_{AC} \otimes (\alpha |0\rangle_B + \beta |1\rangle_B), \quad |\Phi^-\rangle_{AC} \otimes (\alpha |0\rangle_B - \beta |1\rangle_B)$$

$$|\Psi^+\rangle_{AC} \otimes (\alpha |1\rangle_B + \beta |0\rangle_B), \quad |\Psi^-\rangle_{AC} \otimes (\alpha |1\rangle_B - \beta |0\rangle_B)$$

Quantum teleportation: the protocol

4) Communication



Alice tells Bob which of the 4 Bell states she measured using some (classical) binary encoding, e.g.

$$|\Phi^+\rangle_{AC} \mapsto 00$$

$$|\Phi^-\rangle_{AC} \mapsto 01$$

$$|\Psi^+\rangle_{AC} \mapsto 10$$

$$|\Psi^-\rangle_{AC} \mapsto 11$$

Quantum teleportation: the protocol

5) Correction



Bob is now able to correct his qubit in order to recover the initial quantum state $|\phi\rangle$:

- if 00 is received: $|\Phi^+\rangle_{AC} \otimes (\alpha|0\rangle_B + \beta|1\rangle_B)$, nothing to do
- if 01 is received: $|\Phi^-\rangle_{AC} \otimes (\alpha|0\rangle_B - \beta|1\rangle_B)$ he applies Z
- if 10 is received: $|\Psi^+\rangle_{AC} \otimes (\alpha|1\rangle_B + \beta|0\rangle_B)$ he applies X
- if 11 is received: $|\Psi^-\rangle_{AC} \otimes (\alpha|1\rangle_B - \beta|0\rangle_B)$ he applies $Y := ZX$.

Quantum teleportation: discussion

- No violation of the no-cloning theorem: A is not in state $|\phi\rangle$ anymore!
- Intercepting either half of the entangled pair or the classical bits does not give any information about $|\phi\rangle$
- Summary: entangled pair + 1 qubit = entangled pair + 2 bits
- Entanglement swapping: if Alice shares an entangled pair with Bob, he can teleport his qubit to Carol so that now Alice and Carol share an entangled pair
- 1400 km teleportation achieved with the Chinese Micius satellite (2017)